

High Energy Collisions of BHs



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Phys.Rev.Lett.96:071301,2006 (arXiv: hep-th/0512002)

Phys.Rev.Lett.101:161101,2008 (arXiv:0806.1738)

Phys.Rev.Lett., 2009 (arXiv:0907.1252)

New Worlds in Astroparticle Physics, São Tomé (8-10 September)

Very High Energy Collisions



Gravity ends up winning since energy gravitates
The end of particle physics as we know it?

Hoop Conjecture

(Thorne 1972)

- “An imploding object **forms a BH when, and only when**, a circular hoop with a certain critical circumference can be made that encloses the object in all directions. The critical circumference is **2π the Schwarzschild radius** of the object.”

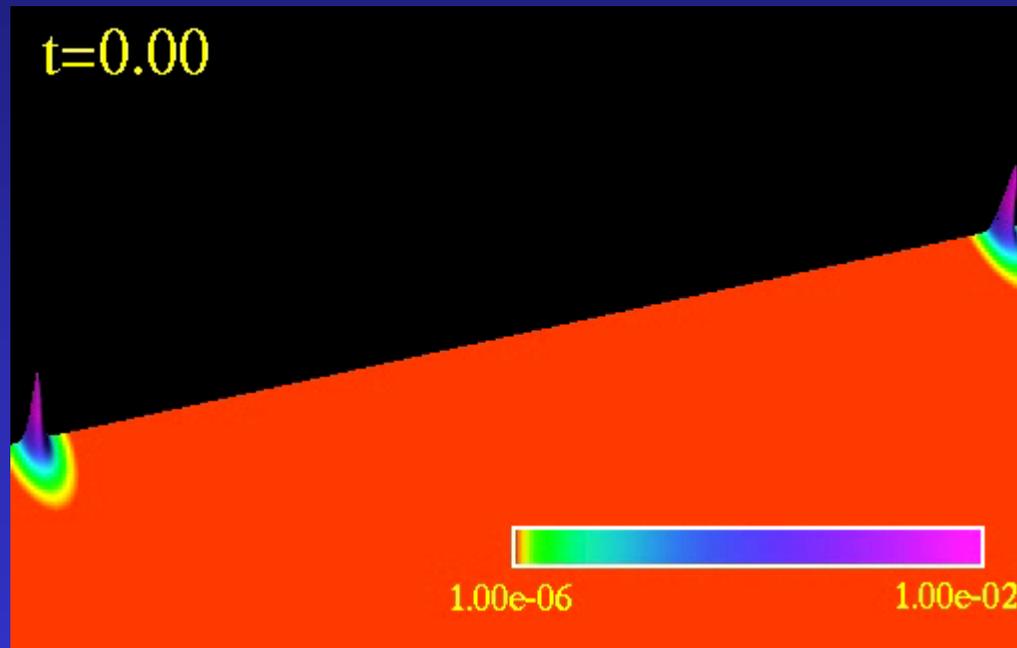
Large amount of energy in small region



This is the hoop
 $R=2GM/c^2$

Hoop Conjecture

(Thorne 1972)



Choptuik & Pretorius, arXiv:0908.1780

Conclusions...

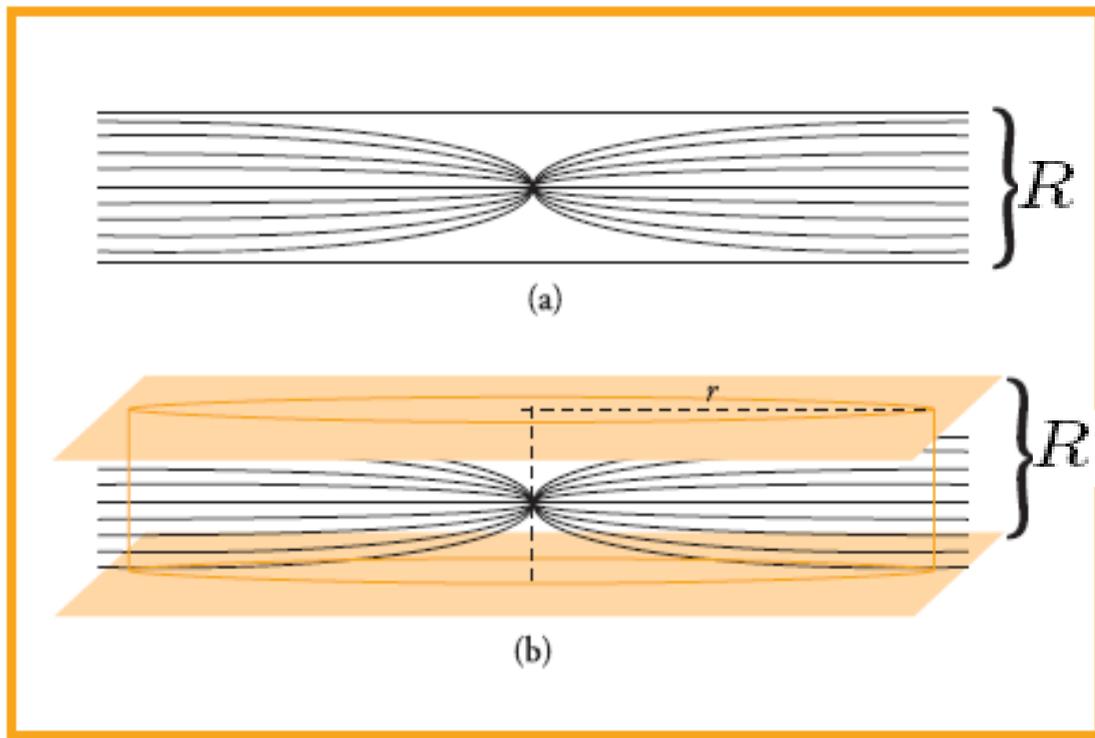
Black holes do form generically in high-energy collisions

Transplanckian scattering is likely to be well described by BH scattering!

The ultra-relativistic 2-body scattering problem is a fascinating theoretical problem within GR, though becomes an issue of more pressing concern if the true Planck scale is within reach of near future experiments...

Brane-New World

- Braneworlds open new possibilities



$$r^{n+2} \times \mathcal{G} = G_{3+n}M$$
$$\mathcal{G} = \frac{G_{3+n}M}{r^{n+2}}$$

$$4\pi r^2 \times R^n \times \mathcal{G} = G_{3+n}M$$
$$\mathcal{G} = \frac{G_{3+n}M}{4\pi R^n r^2}$$
$$\mathcal{G} = \frac{G_3M}{r^2}$$

At large distances, one recovers Newton's gravity!

High-Energy Collisions of BHs

- If Planck scale is small, may be useful for accelerators and Cosmic Ray Experiments!
- Compute production cross section, radiated energies, etc

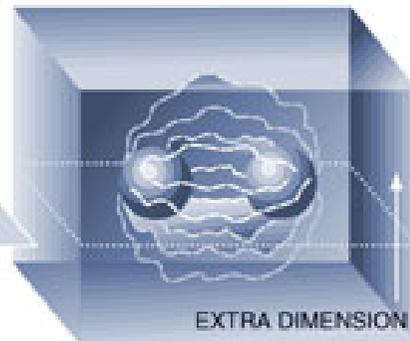
Black Holes on Demand

Scientists are exploring the possibility of producing miniature black holes on demand by smashing particles together. Their plans hinge on the theory that the universe contains more than the three dimensions of everyday life. Here's the idea:

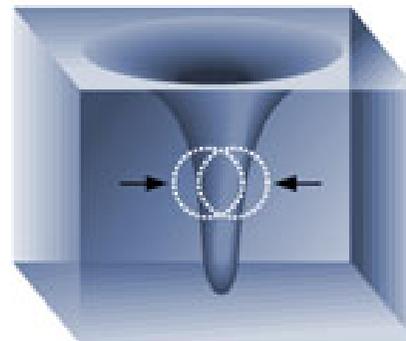
Particles collide in three dimensional space, shown below as a flat plane.



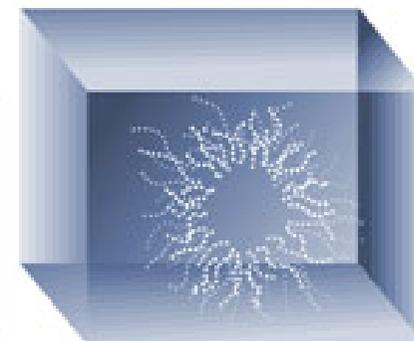
As the particles approach in a particle accelerator, their gravitational attraction increases steadily.



When the particles are extremely close, they may enter space with more dimensions, shown above as a cube.



The extra dimensions would allow gravity to increase more rapidly so a black hole can form.



Such a black hole would immediately evaporate, sending out a unique pattern of radiation.

High-Energy Collisions of BHs

- Cosmic Censorship: do horizons always form?
- Universal Dyson limit on maximum luminosity
(F. Dyson '63)
- Is there critical behavior at the onset of BH formation?
(F. Pretorius '07)
- Test analytical techniques, their predictions and power
(Penrose '74, D'Eath & Payne '93, Eardley & Giddings '02, Cardoso & Lemos '02)
- Holography and HIC
(Amsel et al '08; Gubser et al '08)

**What was known so far
(for head-on collisions)**

I. BH formation: analytics for head-on

- Superpose two Aichelburg-Sexl metrics
 - ✓ Future trapped surface on shock-wave collision
(Penrose '74, Eardley & Giddings '02)
 - ✓ Max. gravitational radiation: $29\% \gamma M_0$



Changes ST from Petrov type D to N

No handle on bounds

No details on wave emission

- Perturb superposed A-S metric, correction: $16\% \gamma M_0$
(D'Eath & Payne '90s)

II. Radiation: ZFL

(Weinberg '64; Smarr '77)

Idea: Infinite acceleration, zero duration process

$$\begin{aligned}g_{\mu\nu} &= \eta_{\mu\nu} + h_{\mu\nu}, \\h_{\mu\nu} &= 4 \int \frac{S_{\mu\nu}(t_r, x')}{\|x' - x\|} d^3x', \\S_{\mu\nu} &= T_{\mu\nu} - \frac{1}{2}\eta_{\mu\nu}T\end{aligned}$$

Take two free particles, changing abruptly at $t=0$

$$\begin{aligned}T^{\mu\nu} &= \sum_{i=1,2} \frac{P_i^\mu P_i^\nu}{E_i} \delta^3(x - v_i t) \theta(-t) \\&\quad + \frac{P_i'^\mu P_i'^\nu}{E_i'} \delta^3(x - v_i' t) \theta(t)\end{aligned}$$

II. Radiation: ZFL

(Weinberg '64; Smarr '77)

$$\frac{dE}{d\omega d\Omega} = \frac{M^2 \gamma^2 v^4}{\pi^2} \frac{\sin^4 \theta}{(1 - v^2 \cos^2 \theta)^2}$$

- i) Flat, large memory ($[h_{t=\infty} - h_{t=-\infty}]^2 \propto dE/d\omega$)
- ii) Radiation isotropic in the limit $\gamma \rightarrow \infty$
- iii) Functional relation $E_{\text{rad}}(\gamma)$
- v) Roughly 65% of maximum possible at $\gamma=3$

With cutoff $M\gamma\omega \approx 0.4$ we get 25% efficiency for conversion of gws



Flat background, linearized approximation!

III. Radiation: PP

(Cardoso & Lemos '02)

Metric=Schwarzschild + small perturbation due to infalling particle

$$h_{\mu\nu} = \begin{bmatrix} H_0(r)f(r) & H_1(r) & 0 & 0 \\ H_1(r) & H_2(r)/f(r) & 0 & 0 \\ 0 & 0 & r^2 K(r) & 0 \\ 0 & 0 & 0 & r^2 K(r) \sin^2 \theta \end{bmatrix} e^{-i\omega t} Y_{l0}(\theta)$$

$$\frac{d^2 \Psi}{dr_*^2} + [\omega^2 - V] \Psi = \left(1 - \frac{2M}{r}\right) S(\omega, r)$$

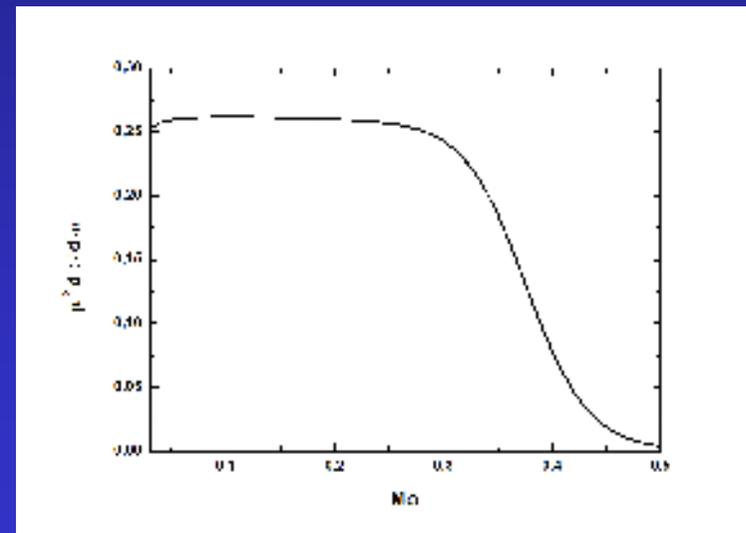
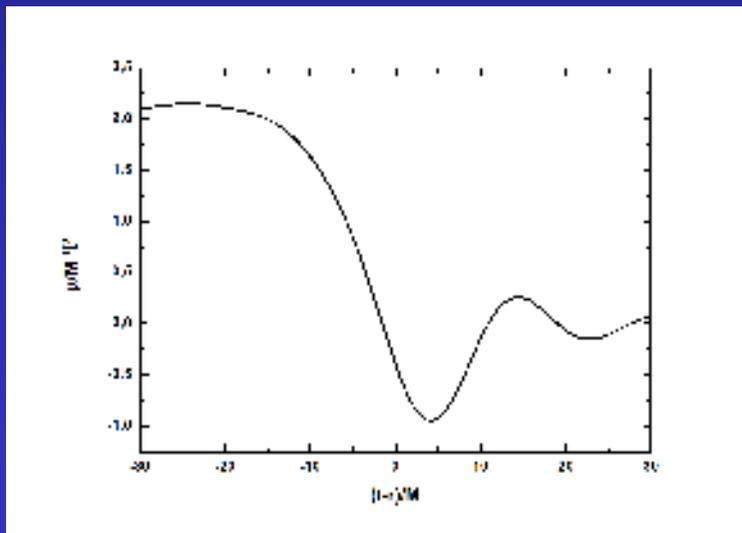
$$\frac{V}{1 - 2M/r} = \frac{2\lambda^2(\lambda + 1)r^3 + 6\lambda^2 r^2 M + 18\lambda r M^2 + 18M^3}{r^3(3M + \lambda r)^2},$$

$$S(\omega, r) = \frac{4i\mu\lambda\sqrt{4l+2}}{\omega(3M + \lambda r)^2} e^{-i\omega r_*}, \quad \lambda \equiv \frac{(l-1)(l+2)}{2},$$

III. Radiation: PP

(Cardoso & Lemos '02)

- i) $dE/d\omega$ flat at sufficiently low ω , large memory
- ii) Functional relation $E_{\text{rad}}(\gamma)$
- iii) Roughly 65% of maximum possible at $\gamma=3$



Linearized, PP approximation!

A very brief history

- Pioneers: Hahn and Lindquist 1960s, Eppley and Smarr 1970s
- Numerical simulations Grand Challenge '90s
Limited accuracy, mild boosts
- Breakthrough: Pretorius 2005, Brownsville, NASA Goddard 2005

Numerical simulations

- LEAN code [Sperhake '07](#)
 - ✓ Based on the Cactus computational toolkit
 - ✓ BSSN formulation (ADM-like, but strongly hyperbolic)
 - ✓ Puncture initial data [Brandt & Brügmann 1996](#)
 - ✓ Elliptic solver: [TwoPunctures](#) [Ansorg 2005](#)
 - ✓ Mesh refinement: [Carpet](#) [Schnetter '04](#)
- Numerically very challenging!
 - ✓ Length scales: $M_{\text{ADM}} \propto \gamma M_0$
 - ✓ Horizon Lorentz-contracted “Pancake”
 - ✓ Mergers extremely violent
 - ✓ Substantial amounts of unphysical “junk” radiation

Gravitational Waves

- Most important result: Emitted gravitational waves (GWs)

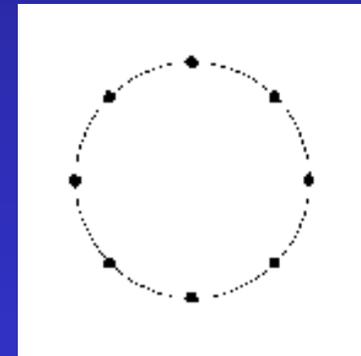
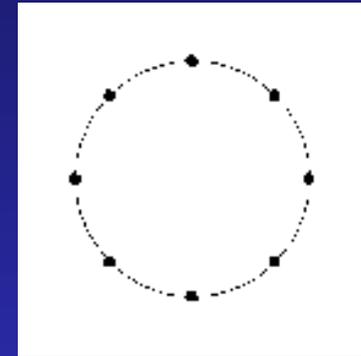
- Newman-Penrose scalar

$$\Psi_4 = C_{\alpha\beta\gamma\delta} n^\alpha m^\beta n^\gamma m^\delta$$

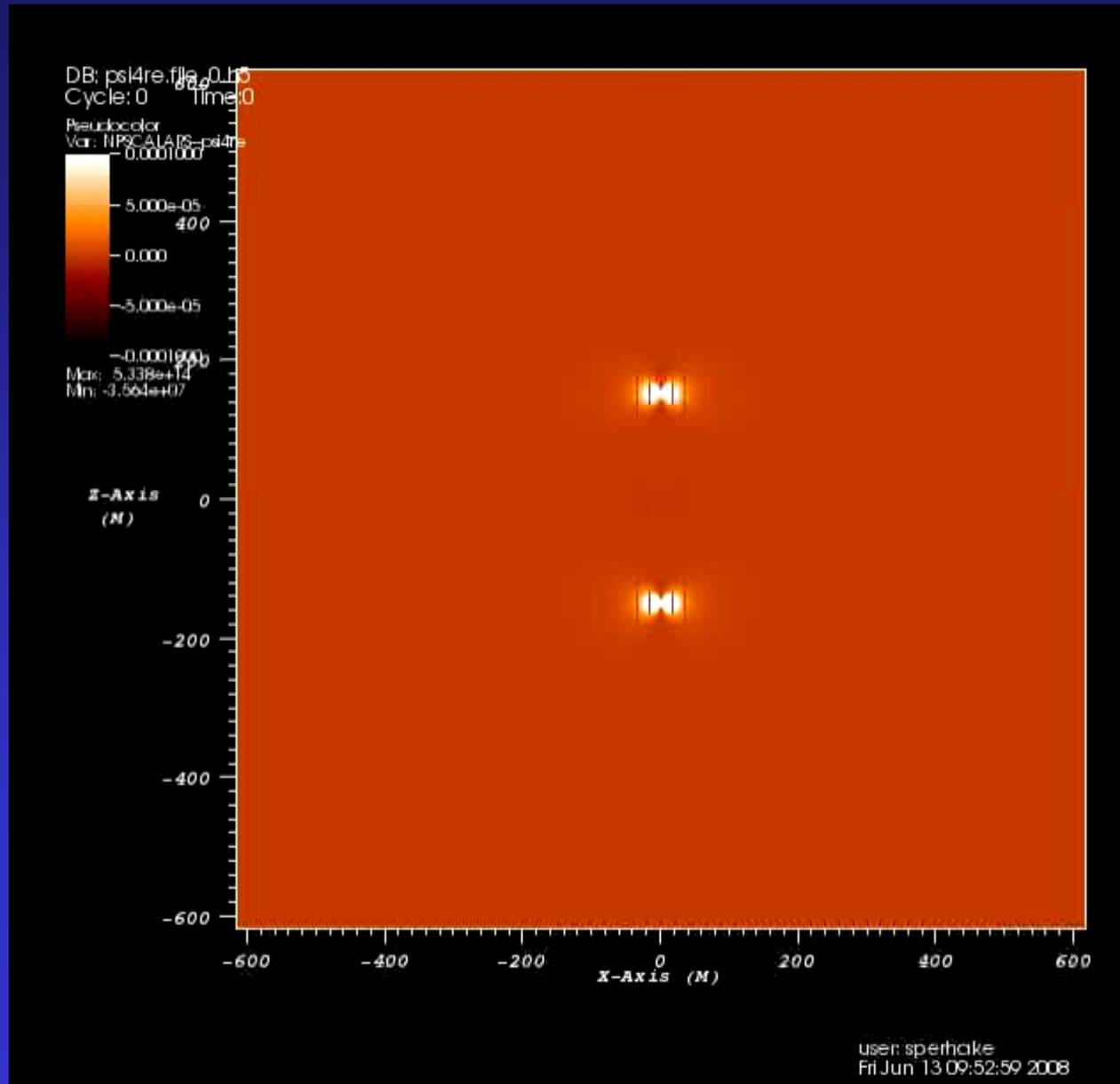
Complex \rightarrow 2 free functions

- GWs allow us to measure

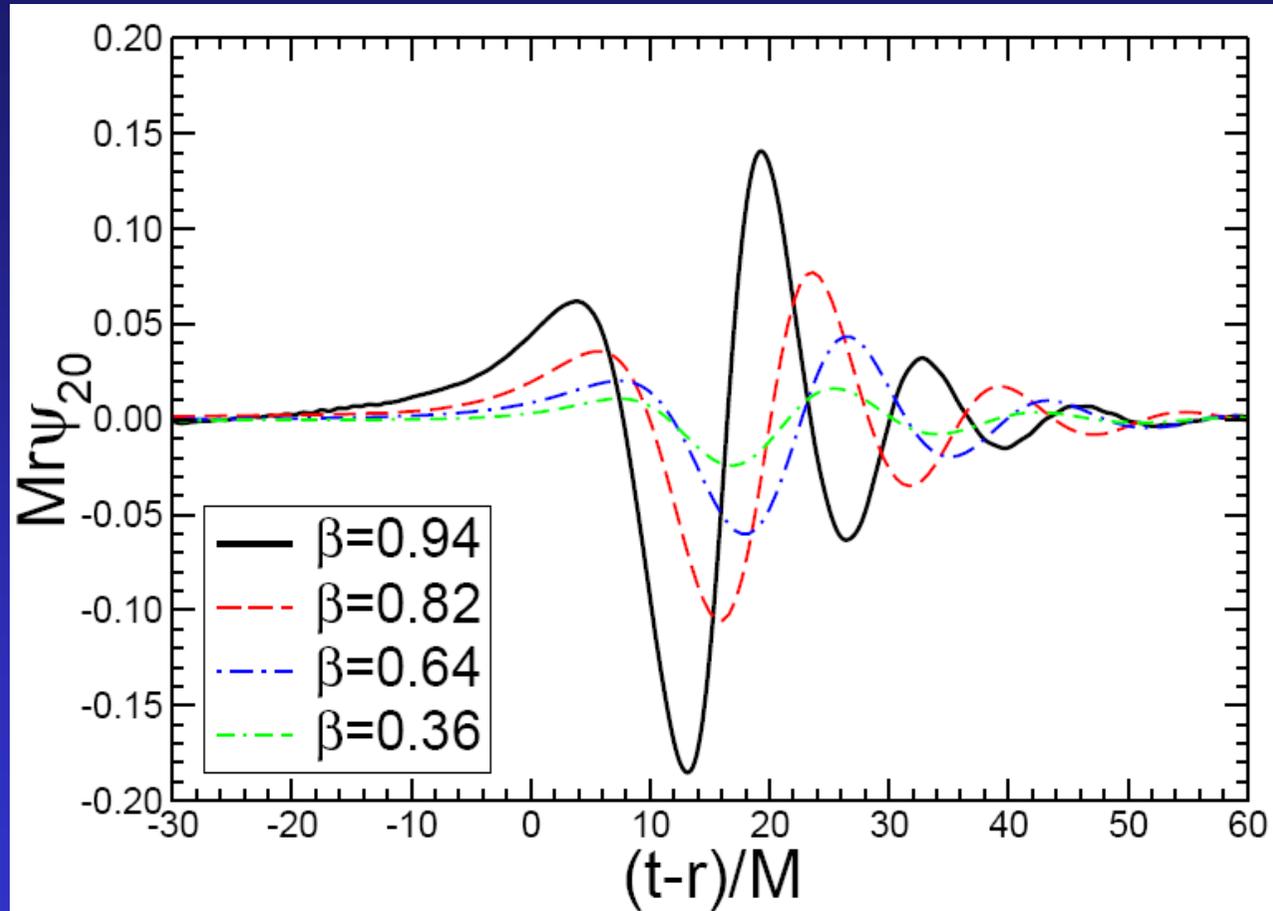
- ✓ Radiated energy E_{rad}
- ✓ Radiated momenta $P_{\text{rad}}, J_{\text{rad}}$
- ✓ Angular dependence of radiation
- ✓ Predicted strain h_+, h_x



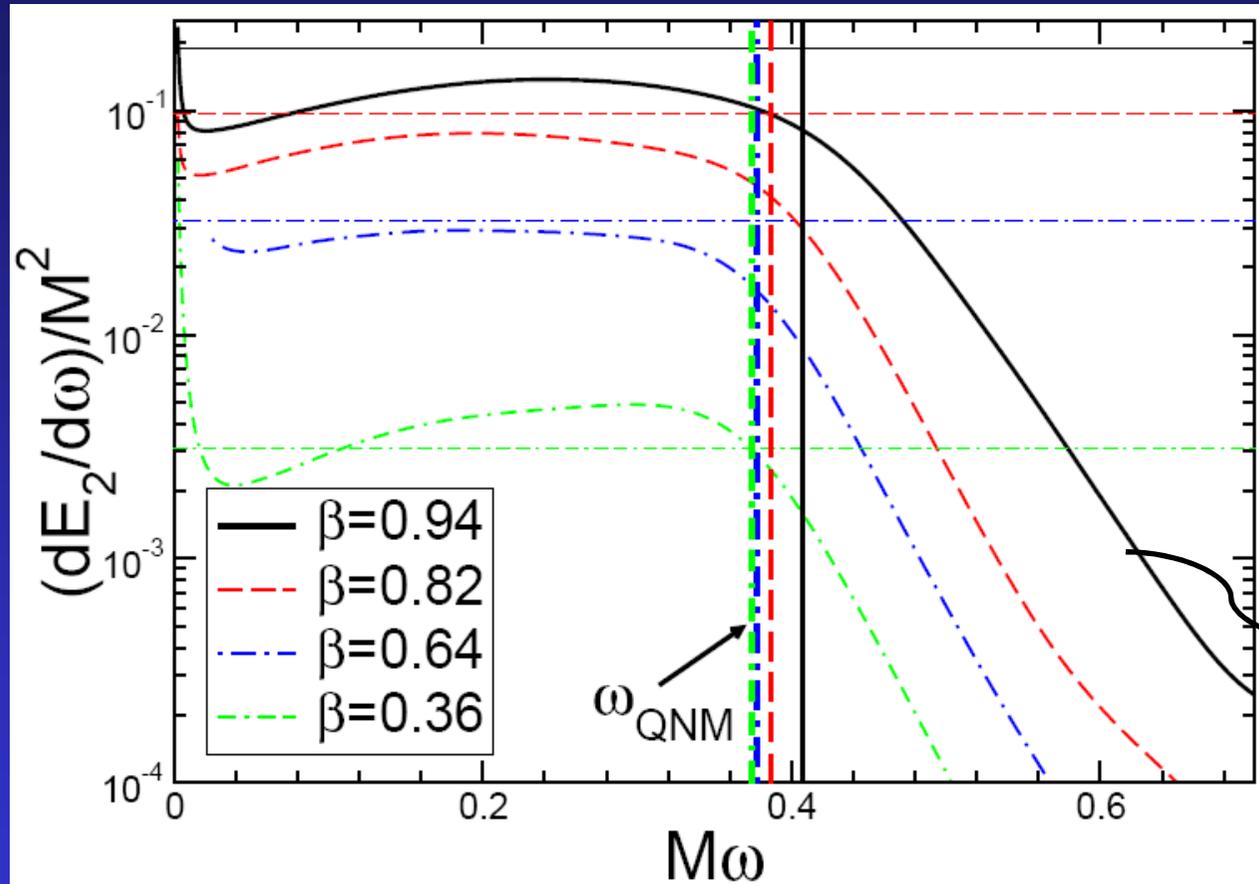
$\beta=0.93$



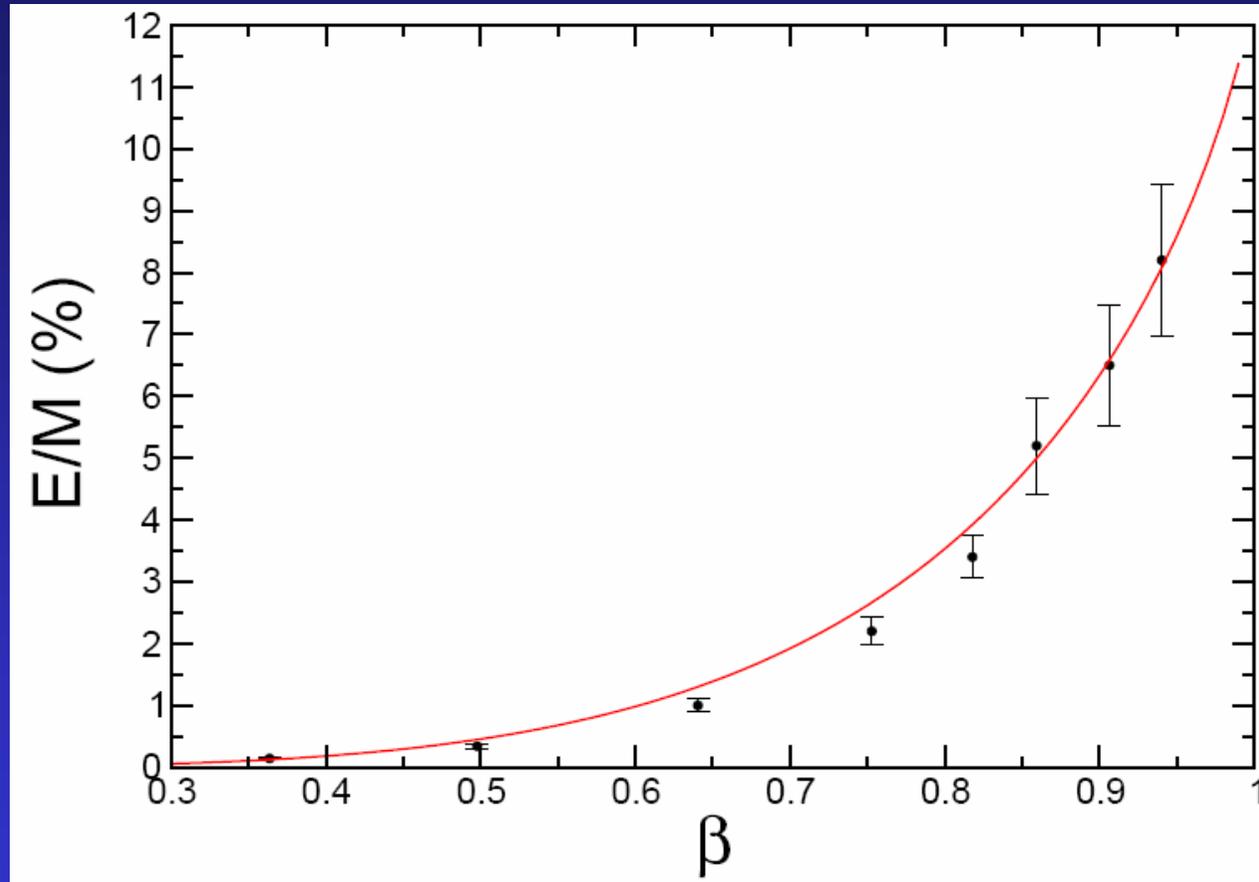
Results



Results



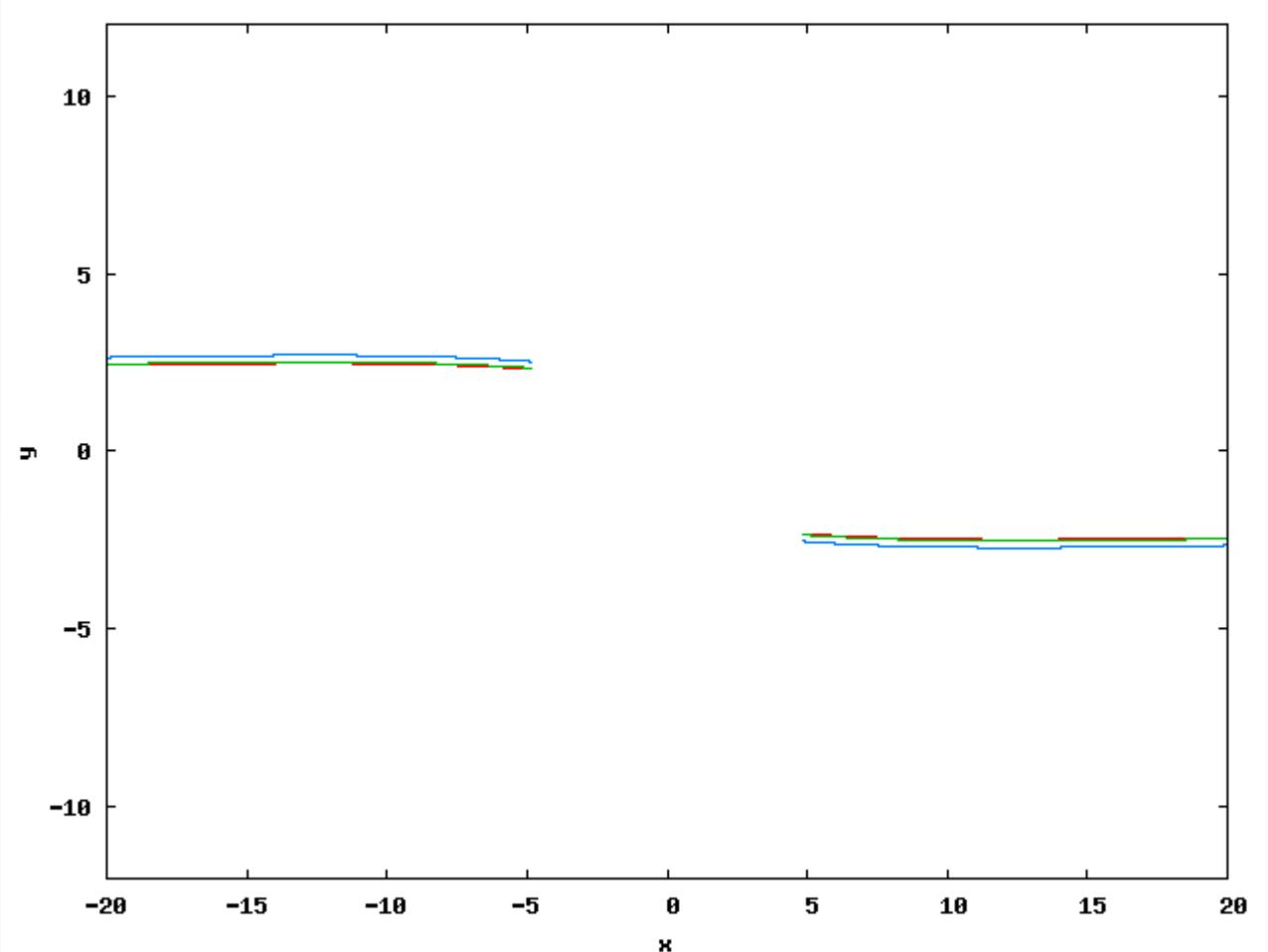
Results



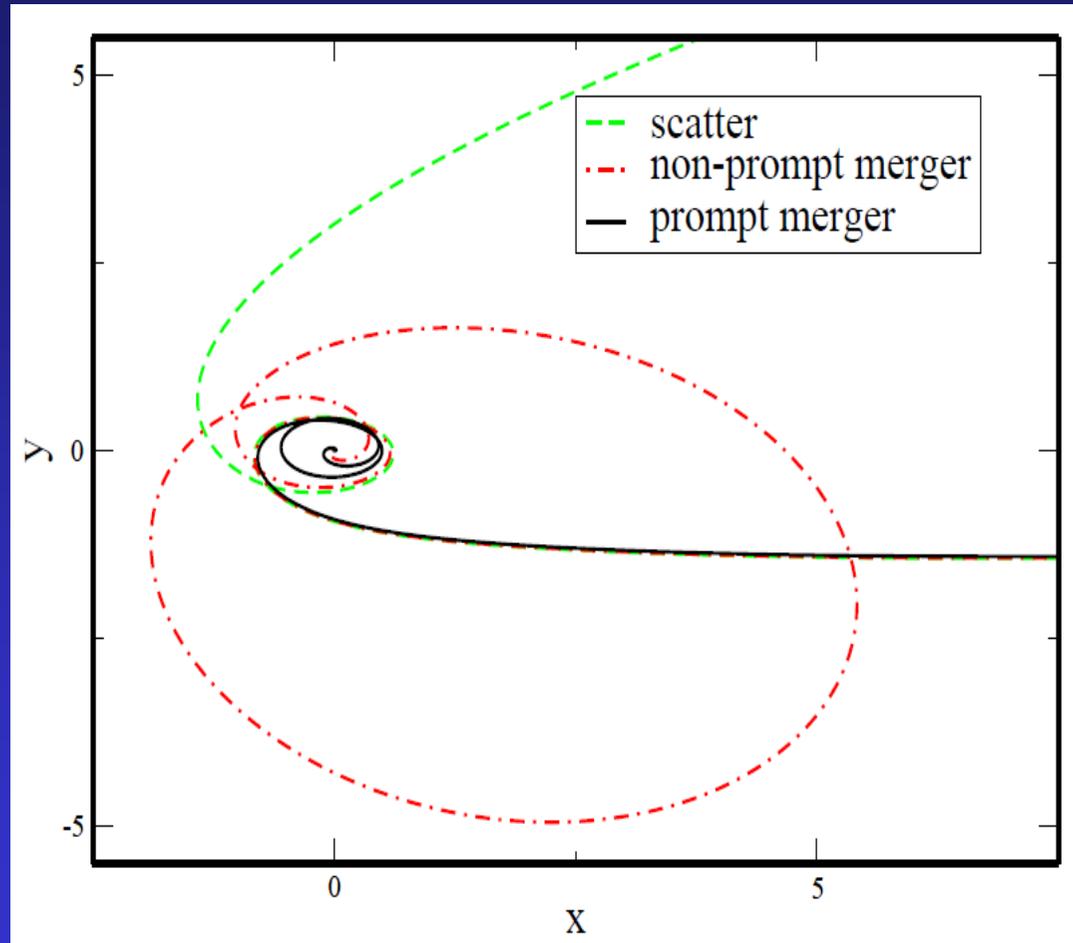
14%

$$\frac{E}{M} = E_\infty \left(\frac{1 + 2\gamma^2}{2\gamma^2} + \frac{(1 - 4\gamma^2) \log(\gamma + \sqrt{\gamma^2 - 1})}{2\gamma^3 \sqrt{\gamma^2 - 1}} \right)$$

Grazing collisions



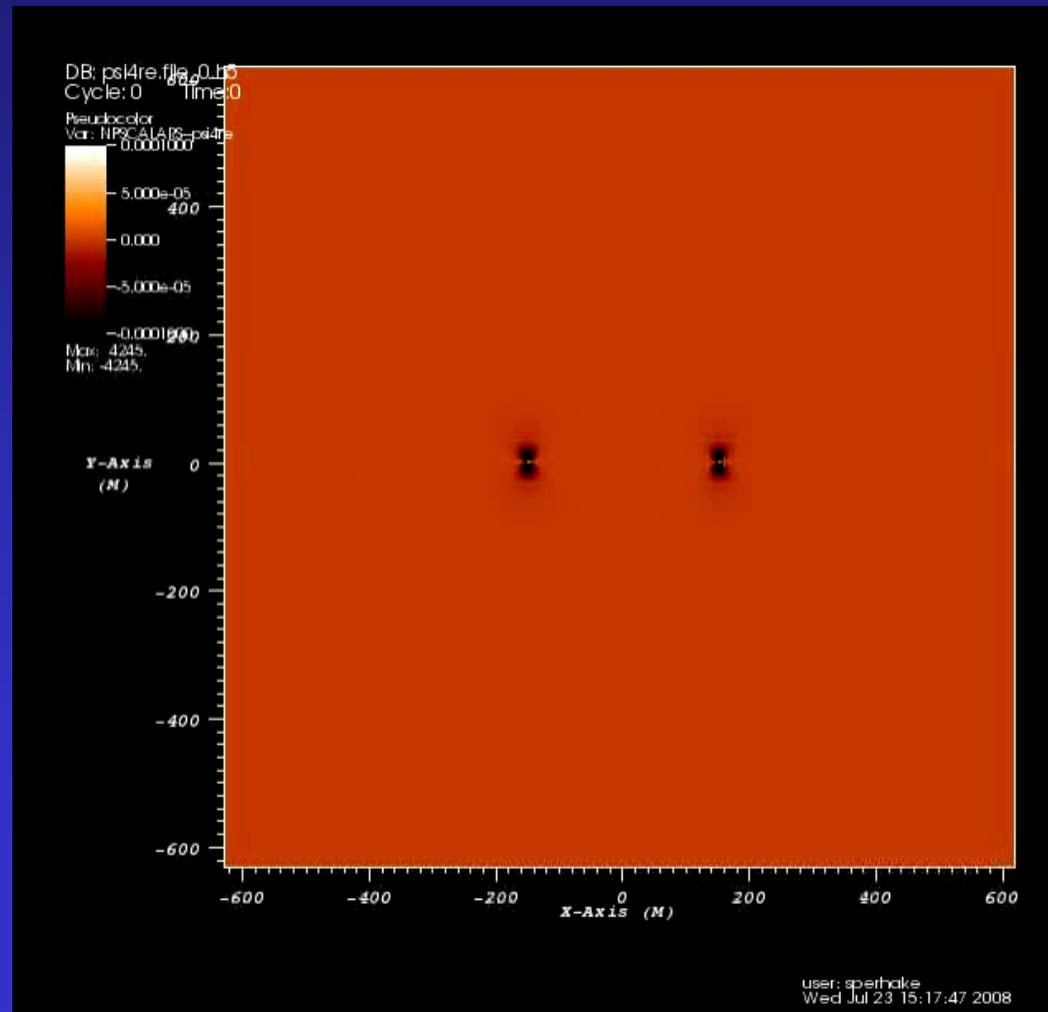
Grazing collisions



Grazing collisions

- ✓ Non head-on: production cross-section

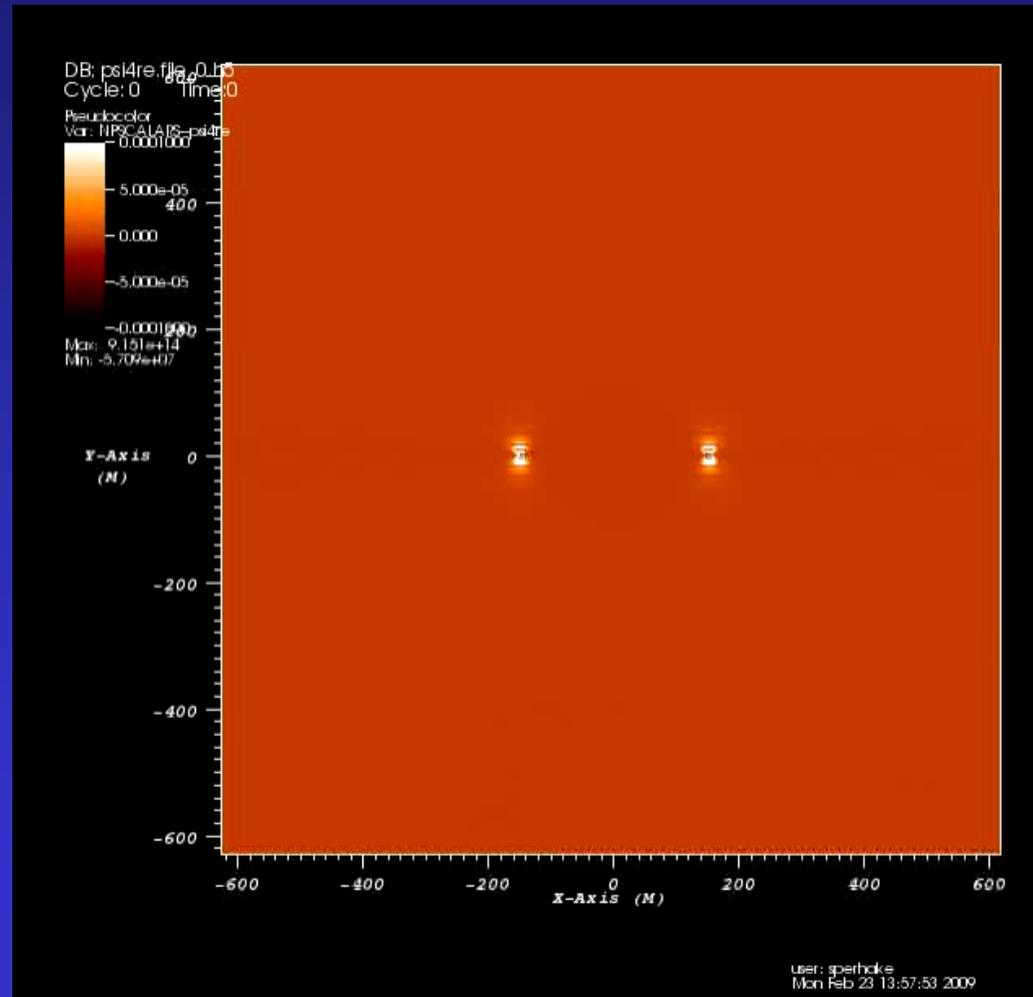
$$\beta=0.75$$



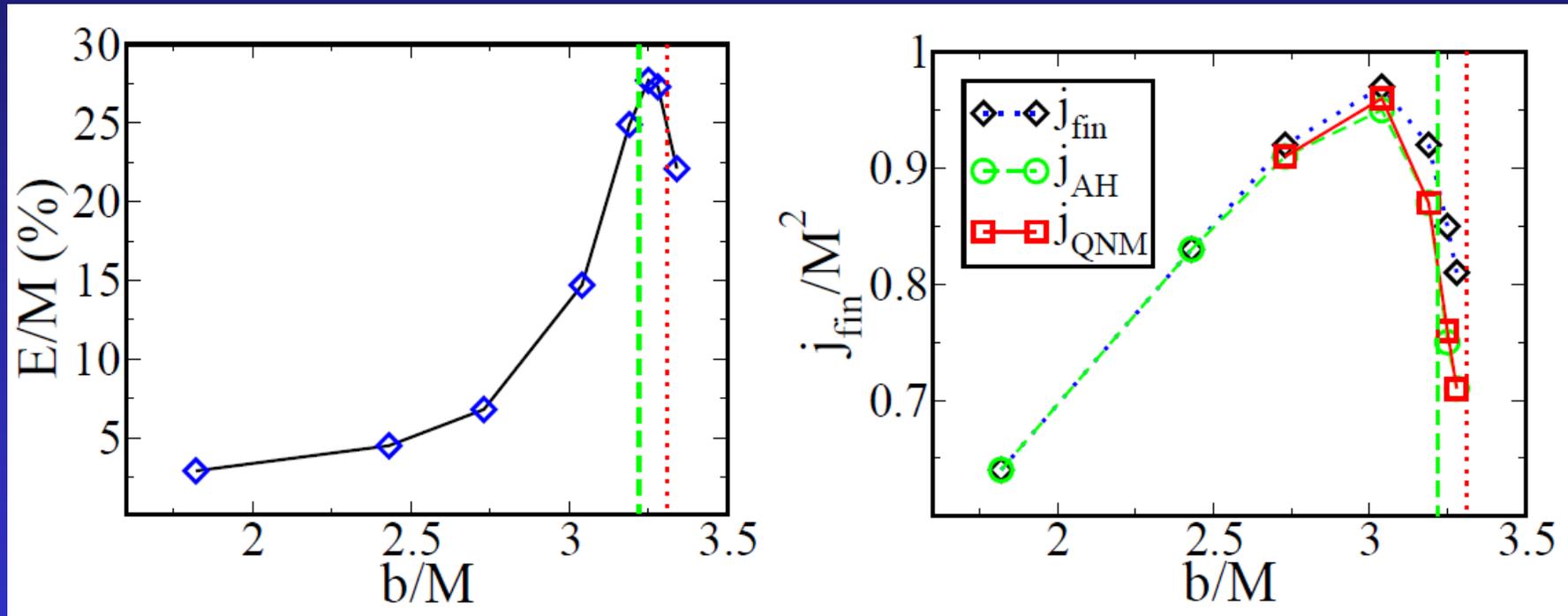
Grazing collisions

- ✓ Non head-on: production cross-section

$$\beta=0.75$$



Grazing collisions



More than 25%CM energy radiated for $v=0.75 c$!

Final BH rapidly spinning

Grazing collisions

- ✓ Non head-on: production cross-section

$$b/M=2.5/v$$

Huge luminosities (almost Dyson limit)

Future work

- ✓ Higher dimensions, other backgrounds
- ✓ Spin, charge
- ✓ Particles with structure (Choptuik & Pretorius)
- ✓ Can the end product be radiation alone?
- ✓ Can the end product be more than one hole?
- ✓ Quantum effects

Thank you

A list of tasks

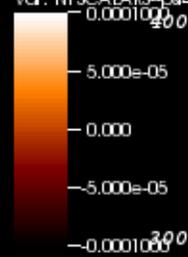
- Target: Predict time evolution of BBH in GR
- Einstein equations:
 - ✓ Cast as evolution system
 - ✓ Choose specific formulation
 - ✓ Discretize for Computer
- Choose coordinate conditions: Gauge
- Fix technical aspects:
 - ✓ Mesh-refinement / spectral domains
 - ✓ Excision
 - ✓ Parallelization
 - ✓ Find large computer
- Construct realistic initial data
- Start evolution and wait...extract physics from data

Gourgoulhon [gr-qc/0703035](https://arxiv.org/abs/gr-qc/0703035)

$$T = 0 M_0$$

DB: psi4re.file.01h5
Cycle: 0 Time: 0

Pseudocolor
Var: NPSCALARS_psi4re



Max: inf
Min: -632.9

Z-Axis
(M)

200

100

0

0

100

200
X-Axis (M)

300

400

$$T = 29.6 M_0$$

DB: psi4re_file_0.h5
Cycle: 4224 Time: 44

Pseudocolor
Var: NPSCALAR5-psi4re-
0.0001000
5.000e-05
0.000
-5.000e-05
-0.0001000
Max: inf
Min: -4.414

Z-Axis
(M)

200

100

0

0

100

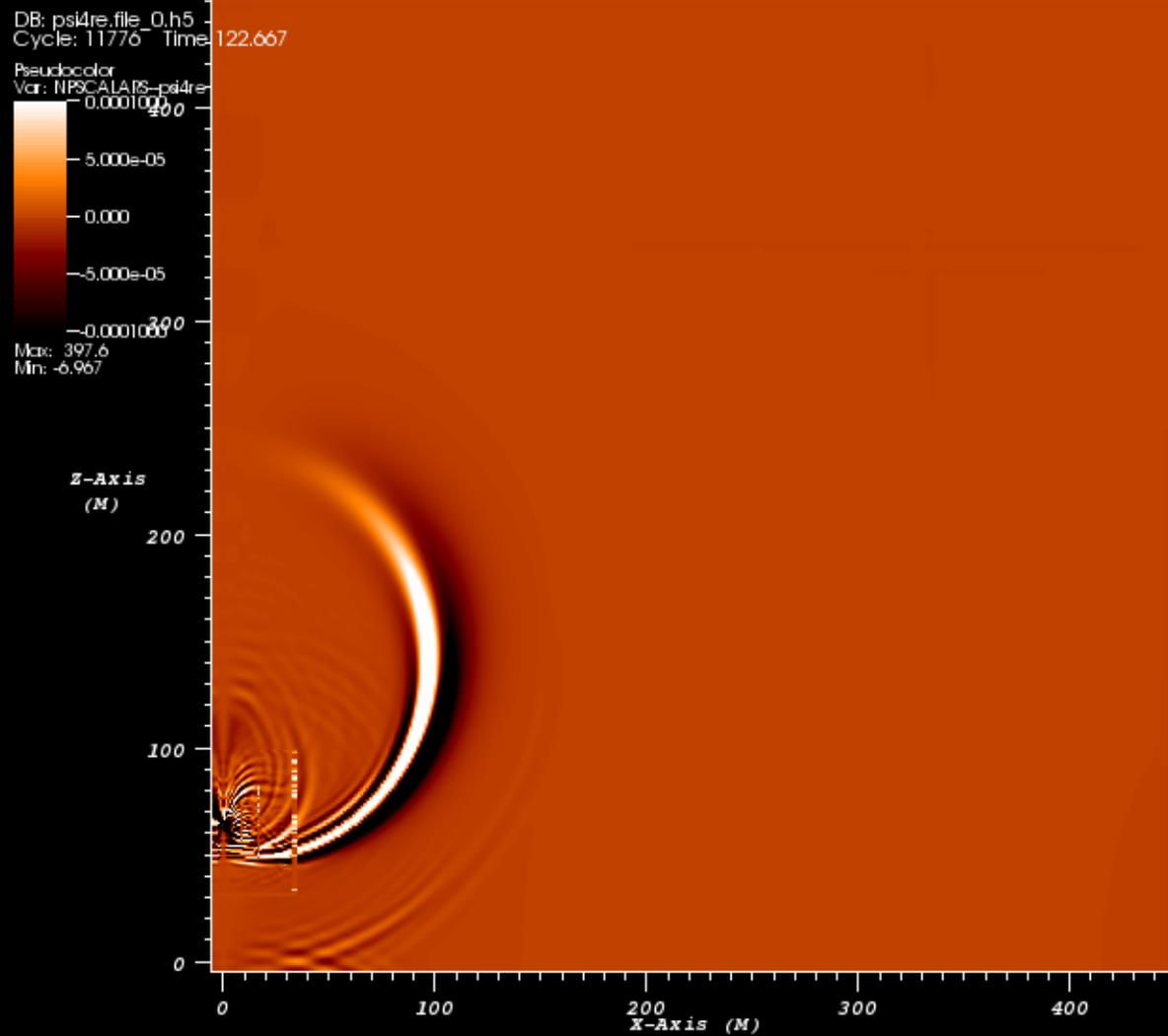
200
X-Axis (M)

300

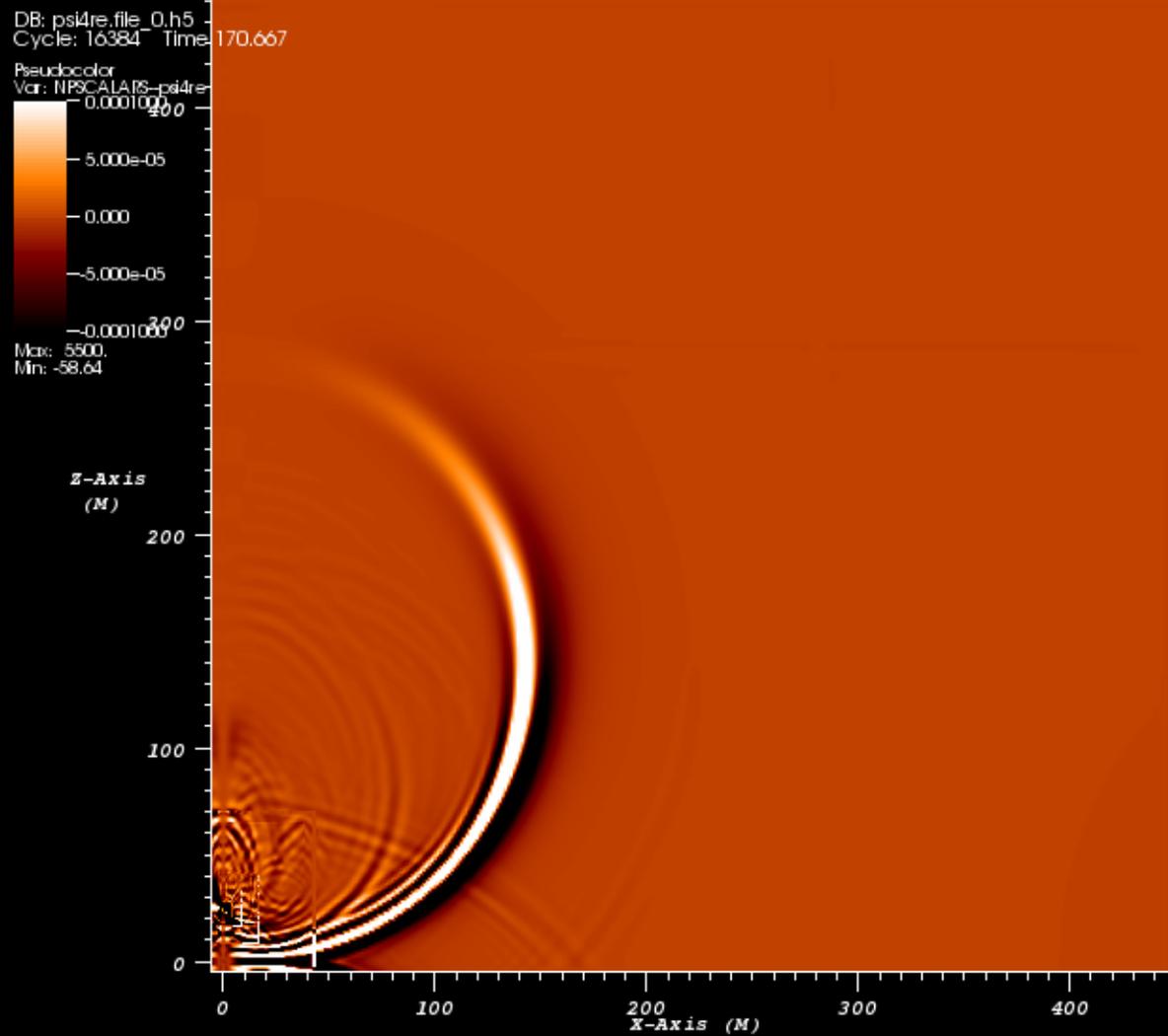
400

user: sperhake
Sun Feb 24 15:41:31 2008

$$T = 82.6 M_0$$



$$T = 114.9 M_0$$



$$T = 137.3 M_0$$

DB: psi4re_file_0.h5
Cycle: 19584 Time: 204

Pseudocolor
Var: NPSCALARS_psi4re
0.0001000
5.000e-05
0.000
-5.000e-05
-0.0001000
Max: 396.5
Min: -2.502

Z-Axis
(M)

200

100

0

0

100

200
X-Axis (M)

300

400

user: sperhake
Mon Feb 25 12:49:59 2008

$$T = 156.1 M_0$$

DB: psi4re_file_0.h5
Cycle: 22272 Time: 232

Pseudocolor
Var: NPSCALAR5-psi4re
0.0001000
5.000e-05
0.000
-5.000e-05
-0.0001000
Max: 396.5
Min: -4887.

Z-Axis
(M)

200

100

0

0

100

200
X-Axis (M)

300

400

user: sperhake
Mon Feb 25 12:58:32 2008

$$T = 201.9 M_0$$

DB: psi4re_file_0.h5
Cycle: 28800 Time: 300

Pseudocolor
Var: NPSCALARS_psi4re
0.0001000
5.000e-05
0.000
-5.000e-05
-0.0001000
Max: 396.5
Min: -614.1

Z-Axis
(M)

200

100

0

0

100

200
X-Axis (M)

300

400

user: sperhake
Mon Feb 25 19:22:24 2008

Grazing collisions

