

Superfluid flow with external localized repulsive potential

VIII BHs Workshop@Lisbon

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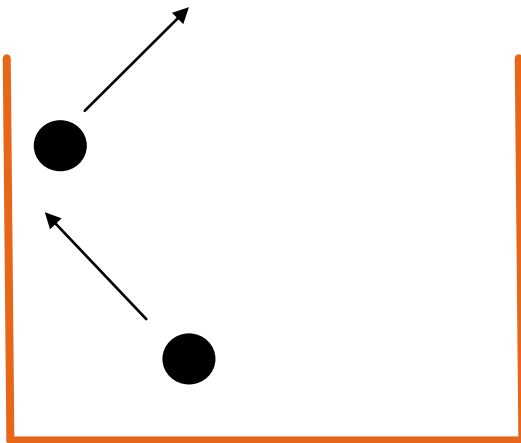
I Introduction

AdS spacetime

AdS/CFT(CMP)



Strongly coupled gauge theory



AdS Boundary



AdS is *non-linearly* unstable

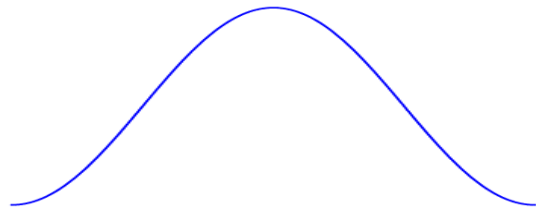


Singularity generically appears from regular initial data!

AdS spacetime \approx Confining BOX

P. Bizon and A. Rostworowski (2011)

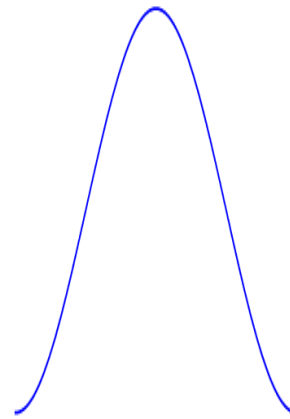
P. Bizon and A. Rostworowski' numerical results



Low frequency modes



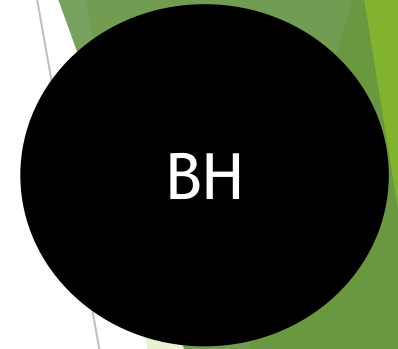
Energy Cascade



High frequency modes



Turbulent behavior



Q. Does any similar turbulent behavior occur under *the black hole background*?

Maybe *Yes!*

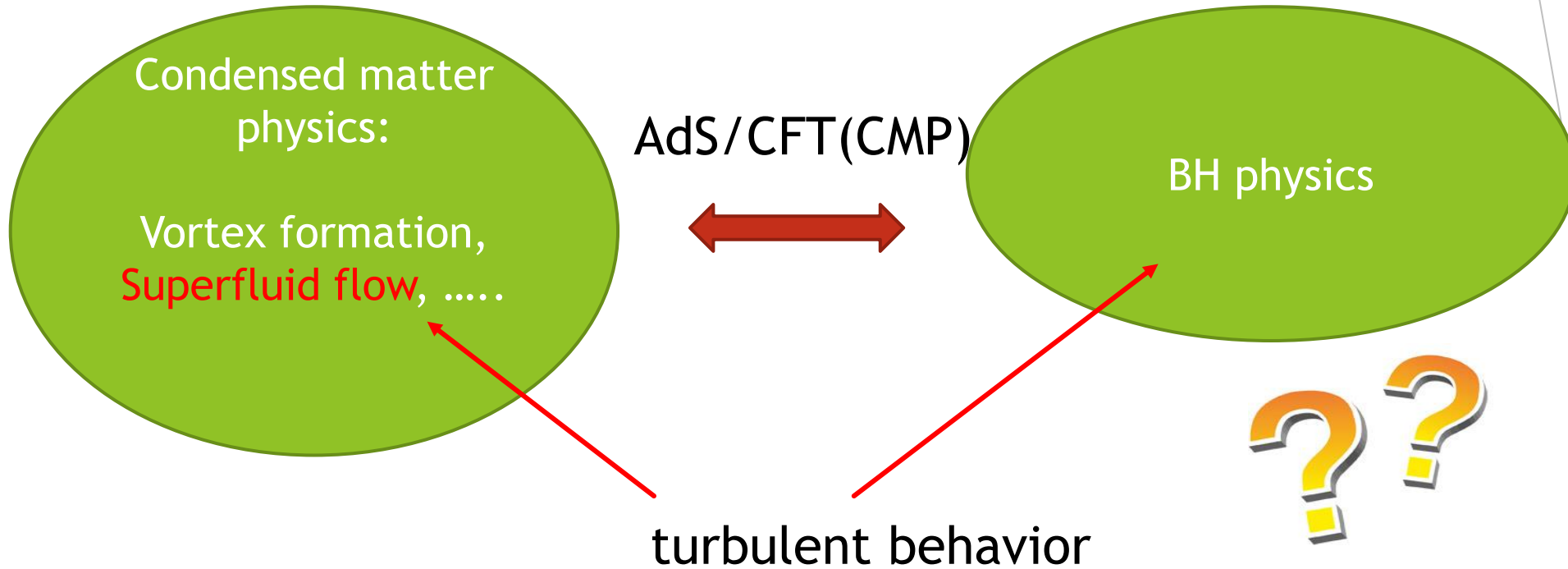
B. E. Niehoff, J. E. Santos, B. Way (2015)

S. R. Green, S. Hollands, A. Ishibashi, R. W. Wald arXiv:1512.02644

Cosmic censorship is violated!?

Superradiant instabilities of Kerr-AdS black holes trigger *energy cascade*

Key question: In a wide class of AdS BHs, does turbulent behavior occur?



Today's talk:

As a first step, we explore how **steady superfluid flow** is broken via AdS/CFT

The model of Bose-Einstein Condensation

1-dim. *Non-linear Schrodinger model*

$$i\partial_t \varphi - i v \partial_x \varphi = -\partial_{xx} \varphi - \varphi + |\varphi|^2 \varphi + U(x) \varphi$$



The limit of application: *Weakly interacting dilute Bose gas system*

Extension to *strongly coupled* gauge theory



AdS/CFT(CMP)

$$\mathcal{L} = -|\nabla\psi - iA\psi|^2 - m^2|\psi|^2 - V(x, u)|\psi|^2 - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$$

u: radial coordinate of AdS

Holographic superconductor model

Hartnoll, Herzog, Horowitz (2008)

Q. How steady superfluid flow is broken?

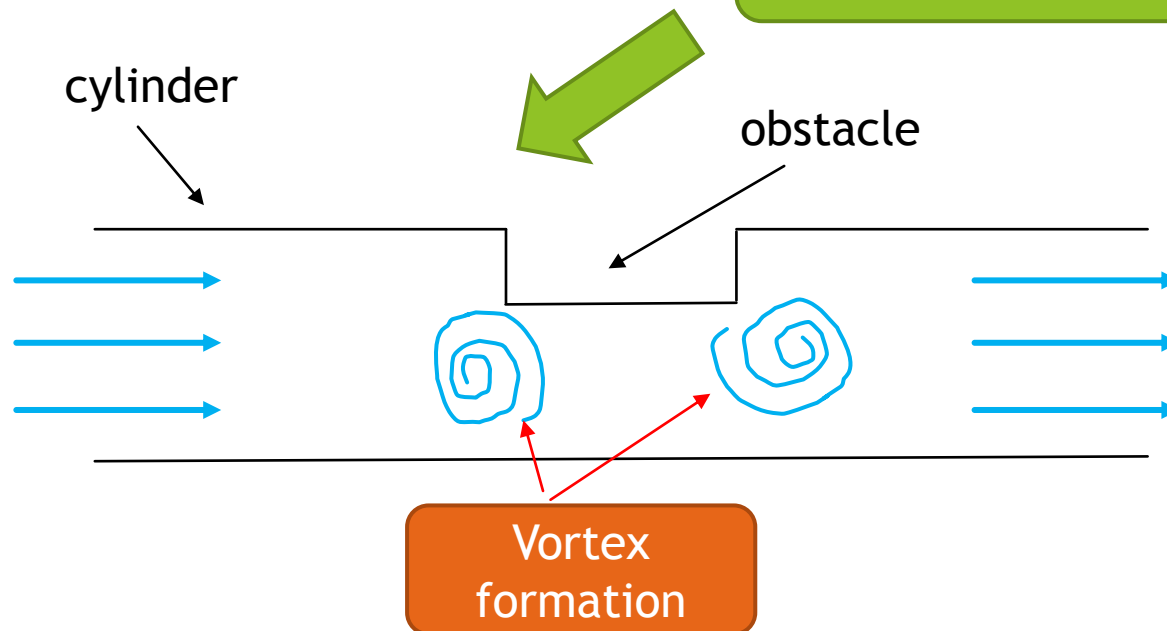
Landau's criterion

$$V > V_L \quad V_L \equiv \text{Min} \varepsilon(p) / p$$

$\varepsilon(p)$: elementary excitation energy

Experimental results: Ex) ^4He

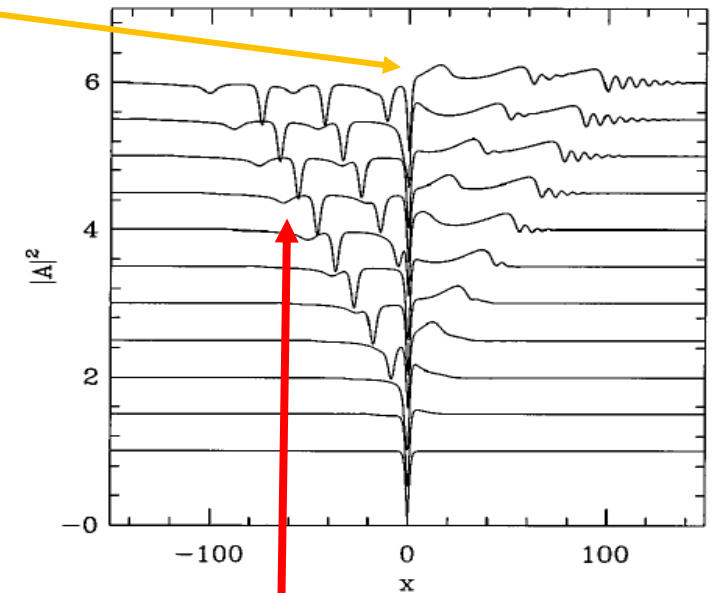
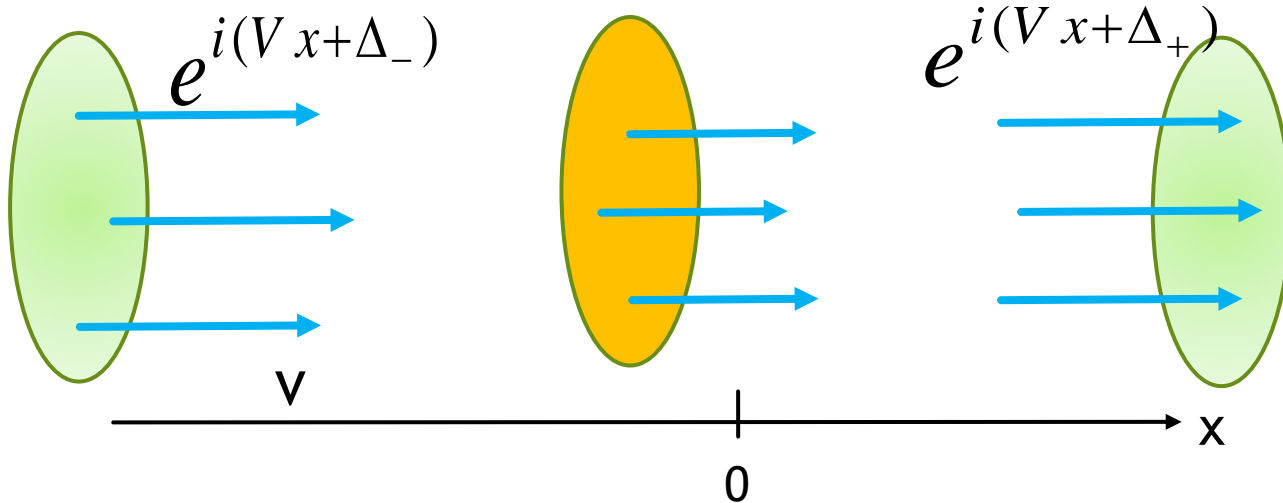
Critical Velocity $V_c < V_L$



Hakim's analysis(1997)

$$i\partial_t \varphi - i v \partial_x \varphi = -\partial_{xx} \varphi - \varphi + |\varphi|^2 \varphi + U(x) \varphi$$

$$U(x) = g \delta(x), \quad g > 0 \quad \text{repulsive potential}$$



- When $g < g_c$ two steady flow solutions appear
- When $g = g_c$ the two solutions coalesce and disappear.
- $g > g_c$ **the gray soliton** solutions are created by the obstacle.

Hakim(1997) PRE55

Set Up of Holographic model

Probe approximation \longrightarrow Sch-AdS5:

$$ds_5^2 = -\frac{1-u^2}{u} dt^2 + \frac{du^2}{4u^2(1-u^2)} + \frac{1}{u}(dx^2 + dy^2 + dz^2).$$

$m^2 = -4$ (Analytic solution can be expected: Herzog(2011))

\longrightarrow $\mathcal{L} = -|\nabla\psi - iA\psi|^2 + 4|\psi|^2 - V(x,u)|\psi|^2 - \frac{1}{4}F_{\mu\nu}F^{\mu\nu}$
 M_μ : Gauge invariant quantity
 $V(x,u) = gu\delta(x)$

Field Eqs.

$$\psi = Re^{i\varphi} \quad M_\mu := A_\mu - \nabla_\mu\varphi$$

$$D^\mu D_\mu\psi + 4\psi - V(x,u)\psi = 0,$$

$$\nabla_\nu F^{\nu\mu} = i[\psi^* D^\mu\psi - \psi(D^\mu\psi)^*],$$

$$D_\mu := \nabla_\mu - iA_\mu$$

$$\nabla^2 R - M^\mu M_\mu R - m^2 R - V(x,u)R = 0,$$



$$\nabla^\mu (M_\mu R^2) = 0$$

Momentum Conservation

$$\nabla_\nu F^{\nu\mu} = 2M^\mu R^2$$

We expand R and M_μ in a series of small parameter ϵ :

$$R = \epsilon R_1(u, x) + \epsilon^3 R_3(u, x) + \dots,$$

$$M_\mu = M_\mu^{(0)}(u, x) + \epsilon^2 M_\mu^{(2)}(u, x) + \dots$$

0th order solution of M_μ

$$M_t^{(0)} = \mu_0(1 - u), \quad \underline{M_x^{(0)} = \zeta(x)}, \quad M_u^{(0)} = 0$$

μ_0 : chemical potential

Velocity of superfluid

AdS/CFT

1th order solution of R_1

$$R_1 = \rho(u)\zeta(x)$$

Separation of variables

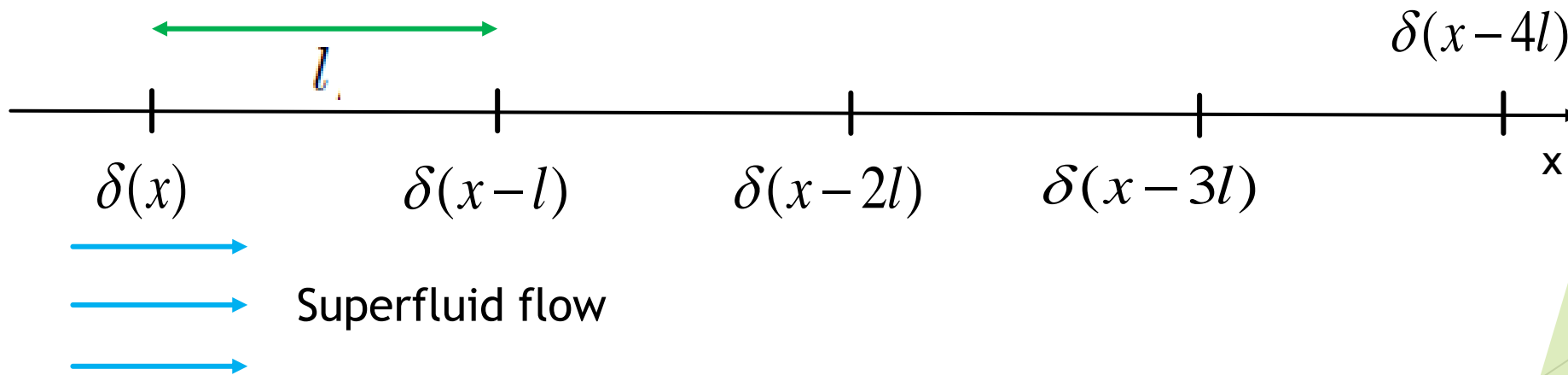
$$\zeta^2 \xi = c_2.$$

Separation constant

$$\partial_x^2 \zeta - \xi^2 \zeta - g\delta(x)\zeta = -\kappa^2 \zeta$$

$$4u^2(1-u^2)\rho'' - 4u(1+u^2)\rho' + \left[4 + \frac{\mu_0^2 u(1-u)}{1+u} - \kappa^2 u \right] \rho = 0.$$

We assume a *periodic potential* V:



Bloch's theorem



ζ must be a periodic function of x

Solution: $0 \leq x \leq l$

$$\zeta^2 = \frac{1 - \alpha^2}{2} \cos 2\kappa(x - l/2) + \frac{1 + \alpha^2}{2}$$

$$c_2 = \alpha\kappa$$

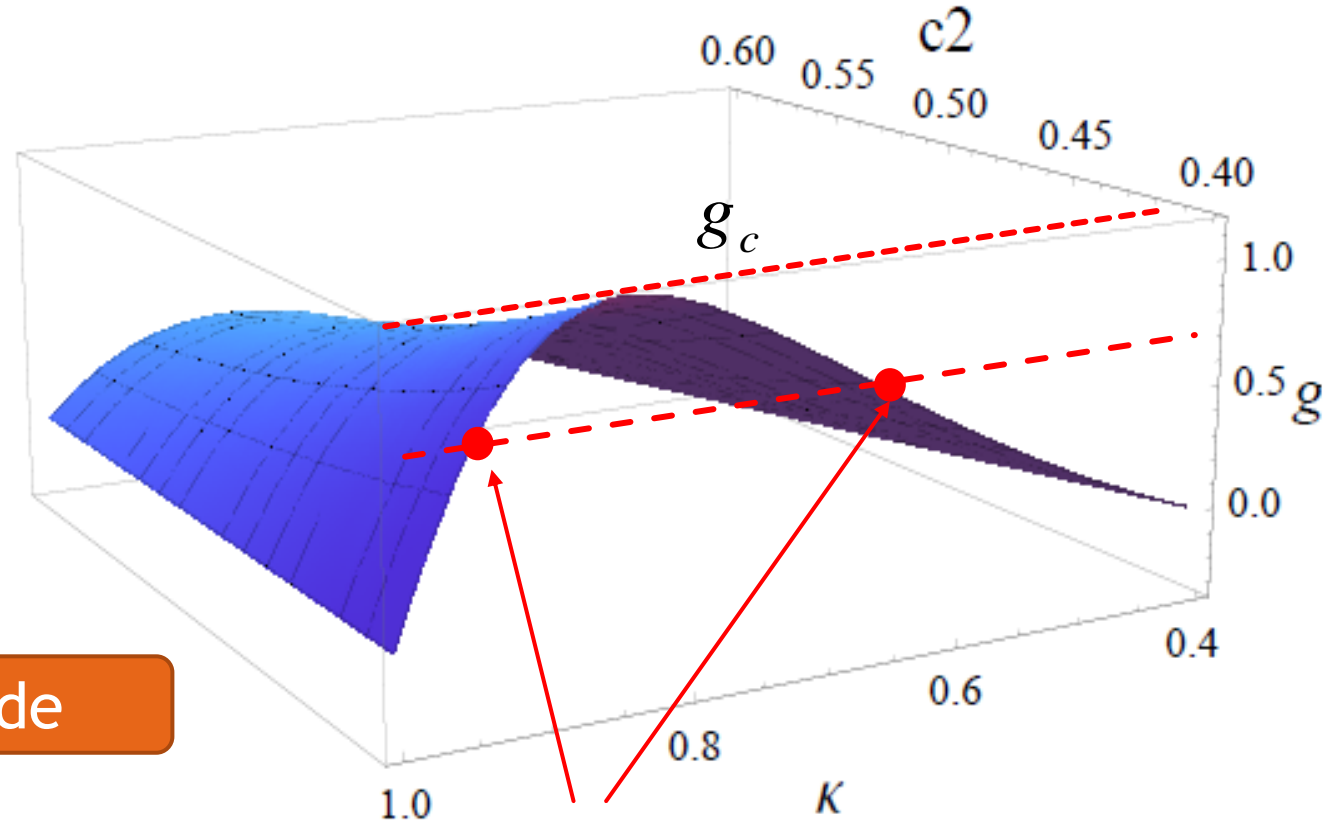
Asymptotic Condition for ρ

$$\rho \rightarrow O(u)$$

Normalizable mode

$$\rho = \frac{u}{1+u} - \kappa^2 \frac{u \ln(1+u)}{4(1+u)} + O(\kappa^4),$$

$$\mu_0 = 2 + \frac{1}{2}\kappa^2 + O(\kappa^4).$$



Two solution coexist

Holographic extension to Hakim's result

2th order solution of $O(\epsilon^2)$

Eqs. for $M_x^{(2)}$, $M_t^{(2)}$

$$(1 - u^2)M_x''^{(2)} - 2uM_x'^{(2)} = \frac{c_2\rho^2}{2u^2}$$

$$4u\partial_u^2 M_t^{(2)} + \frac{1}{1-u^2}\partial_x^2 M_t^{(2)} = \frac{2\mu_0}{u(1+u)}\rho^2\zeta^2$$

$$M_x^{(2)} = \frac{c_2}{32(1+u)} \times$$

polygamma function

$$\left[8(1 - \kappa^2) + \kappa^2 \{ (2 \ln 2)(1 + u) \ln(1 - u) - \ln(1 + u)(4 + 2(1 + u) \ln(1 - u) - (1 + u) \ln(1 + u)) \} \right. \\ \left. - 2\kappa^2(1 + u) \text{Li}_2\left(\frac{1 + u}{2}\right) \right] - \frac{c_2}{192} \{ 48 - \kappa^2(48 + \pi^2 - 6(\ln 2)^2) \} + O(\kappa^4).$$

AdS/CFT



Current of superfluid flow

$$\langle J_x \rangle \sim M_x'^{(2)}(0) = -\frac{c_2}{8}(2 - \kappa^2(1 - \ln 2)).$$

$$M_t^{(2)}(u, x) = \eta_{t0}(u) + \eta_{t1}(u) \cos 2\kappa(x - l/2)$$



What happens near $g = g_c$?

Y. Kato, S. Watabe (2010) PRL105



The analysis of non-linear Schrodinger Eq.

Dynamical density fluctuations (DDFs) is important near the critical velocity



Spectral function of the local density I

$$I(x, \omega) = \sum_l |\langle l | \psi^+(x) \psi(x) | g \rangle|^2 \delta(\omega - E_l + E_g)$$

g: ground state
l: excited state

$$I(x, \omega) = \omega^{d-2} F(x, \omega | g - g_c |^{-\frac{1}{2}})$$

d=1



$$I(x, \omega) \rightarrow \infty \quad \text{for} \quad \omega \rightarrow |g - g_c|^{\frac{1}{2}}$$

Holographic calculation of $I(x, \omega)$



Construction of retarded Green function G_R of perturbed ψ in the bulk

$$D^2 \delta\psi - V(x, u) \delta\psi = 0 \quad \longrightarrow \quad D^2 G_R - V(x, u) G_R = \delta(t - t', x - x', u - u')$$

Separation of variables $\delta\psi = \chi_\kappa(x) \times R(u; \kappa)$

$$(D_x^2 - V(x)) \chi_\kappa = -\kappa^2 \chi_\kappa \quad \sum_\kappa \chi_\kappa^+(x) \chi_\kappa(x') = \delta(x - x')$$

$$\delta\psi \cong \alpha_+(\omega, \kappa) u^{\lambda_+} + \alpha_-(\omega, \kappa) u^{\lambda_-} \quad u \rightarrow 0$$

Normalizable mode

$$I(\omega, x) \approx \text{Im} \left[\sum_\kappa \frac{\alpha_+(\omega, \kappa)}{\alpha_-(\omega, \kappa)} \chi_\kappa^+(x) \chi_\kappa(x) \right]$$

Summary

- We investigated holographic superfluid flow as a first step to analyze *turbulence* of rotating AdS BH
- We obtained *analytic solutions with external repulsive potential* and reproduced Hakim's results in holographic setting



Future plan

- Derivation of spectral function $I(\omega, x)$ in the holographic setting
- Construction of *gray soliton solution* in rotating hairy AdS solutions

Thank you!